

McLean seminar

2024 7/12 Fri.

Riku Sato (Waseda Univ.)

5. Instrumentation and detectors

5.3 POLARIMETERS

Measure polarization properties of light (fraction polarized, direction of vibration, handedness of rotation).

→ Provides unique geometric information not obtainable from intensity alone.

Targets:

- Reflection nebulae
- Synchrotron emission from supernova remnants
- Cyclotron emission from magnetic white-dwarf systems

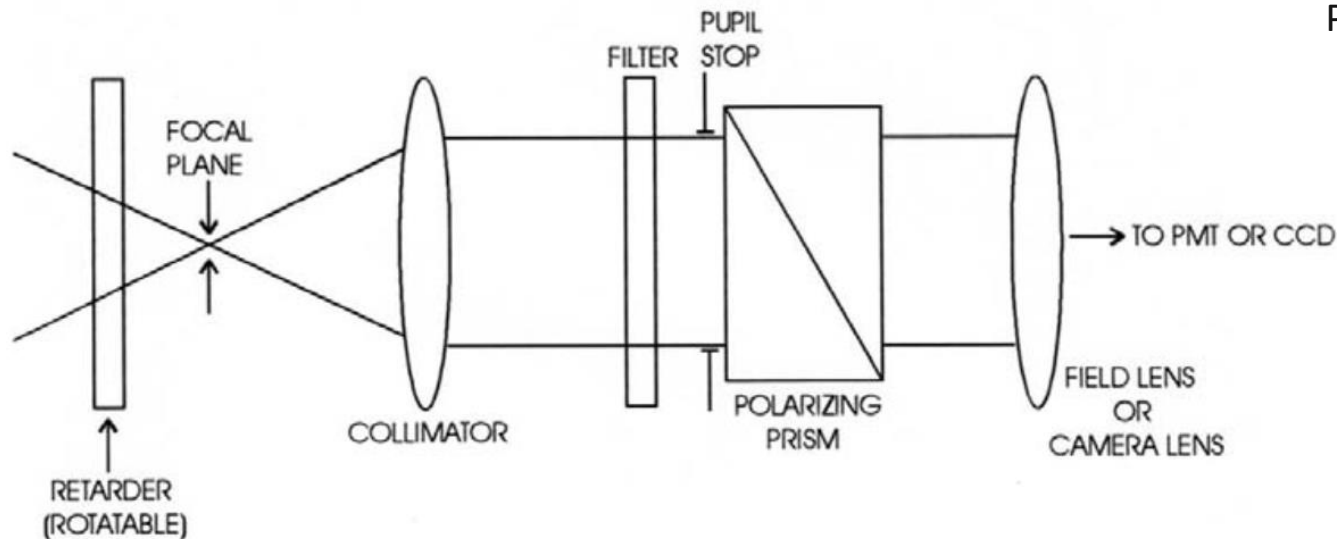
Polarization Information is converted into brightness modulations.

Polarimeter Types:

- Photopolarimeter
- Imaging Polarimeter
- Spectropolarimeter

5.3.1 Modulators and polarizers

- Modulator: Rotatable plate (e.g. quartz and magnesium fluoride) or Pockels cell. Introduces birefringence, sensitive to light wave orientation.
- Polarizer: Typically glass components (e.g. Glan-Thompson or Wollaston prism). Allows controlled brightness variation based on polarization.
- Light from telescope passes through polarimeter (modulator + polarizer) then to detector.
- Modulation of brightness occurs due to polarized light.
- Detector records brightness variations.



Polarizer: A specific linear polarization can only pass

Rotator: Change the angle of polarization

Retarder: Make phase shift in the beam

5.3.2 The Stokes parameters

Stokes Parameters:

Linear Polarization: Intensity (I), degree (p), direction (θ).

Circular Polarization: Intensity (I), degree (q), handedness of the rotation of the electric vector (+ or -).

A more convenient way \rightarrow use four Stokes Parameters: I, Q, U, V.

I: Total Intensity, Q, U: linear polarization, V: Circular polarization

degree of polarization p: linear, q: circular

$$p = \frac{[Q^2 + U^2]^{1/2}}{I}, \quad q = \pm \frac{V}{I} \quad (5.30)$$

and the direction of vibration of the linearly polarized part is given by

$$\tan 2\theta = \frac{U}{Q} \quad (5.31) \quad \text{where}$$

and it follows that

$$\left. \begin{aligned} Q &= Ip \cos 2\theta \\ U &= Ip \sin 2\theta \\ V &= Iq \end{aligned} \right\} \quad (5.32)$$

$$I' = \frac{1}{2} [I \pm Q(G + H \cos 4\psi) \pm UH \sin 4\psi \mp V \sin \tau \sin 2\psi] \quad (5.33)$$

$$G = \frac{1}{2}(1 + \cos \tau), \quad H = \frac{1}{2}(1 - \cos \tau), \quad \tau = \frac{2\pi}{\lambda} \delta \quad (5.34)$$

5.3.2 The Stokes parameters

A more convenient way → use four Stokes Parameters: I, Q, U, V.

I: Total Intensity, Q, U: linear polarization, V: Circular polarization

(I, Q, U, V)

= (1, 1, 0, 0) : linear (horizontal)

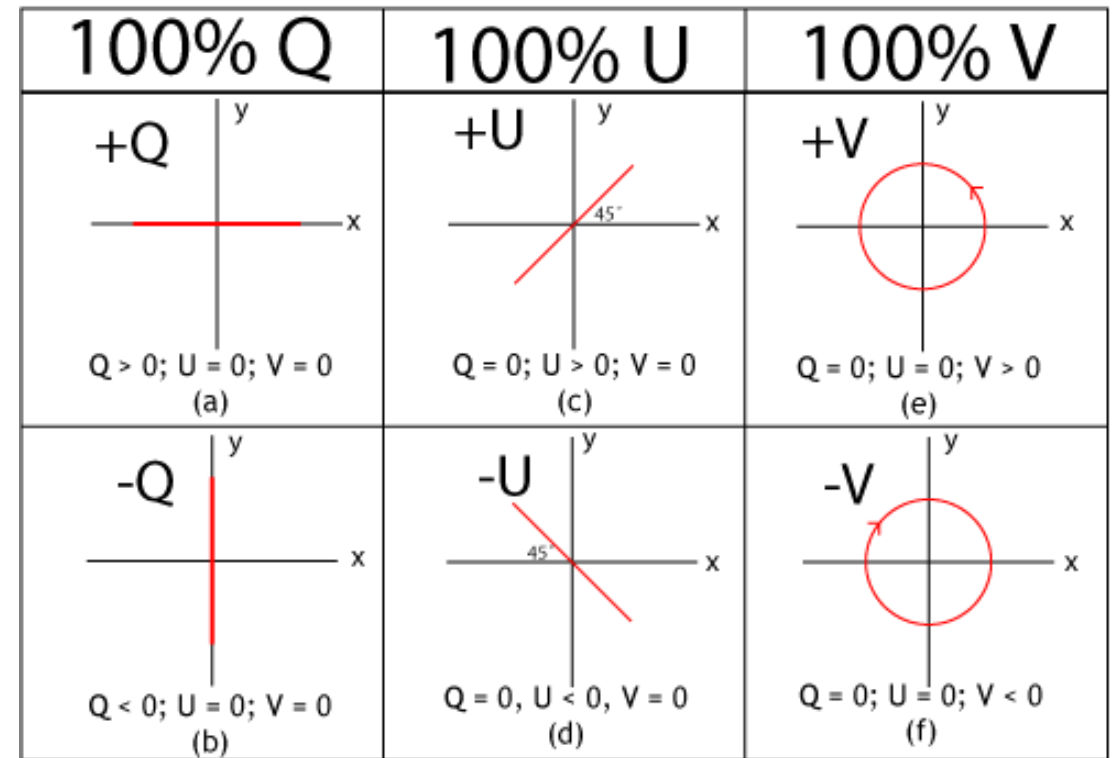
= (1, -1, 0, 0) : linear (vertical)

= (1, 0, 1, 0) : linear (+45°)

= (1, 0, -1, 0) : linear (-45°)

= (1, 0, 0, -1) : circular (clockwise)

= (1, 0, 0, 1) : circular (counterclockwise)



5.3.2 The Stokes parameters

A more convenient way → use four Stokes Parameters: I, Q, U, V.

I: Total Intensity, Q, U: linear polarization, V: Circular polarization

(1) *The Quarter-Wave Retarder*: $\delta = \lambda/4$, $\tau = 90^\circ$, $G = H = \frac{1}{2}$ which gives

$$I' = \frac{1}{2} [I \pm \frac{1}{2} Q \cos 4\psi \pm \frac{1}{2} U \sin 4\psi \mp V \sin 2\psi] \quad (5.35)$$

Solar magnetographs must determine the circular component, and Method (1) is the basis for those instruments.

(2) *The Half-Wave Retarder*: $\delta = \lambda/2$, $\tau = 180^\circ$, $G = 0$, $H = 1$ which gives

$$I' = \frac{1}{2} [I \pm Q \cos 4\psi \pm U \sin 4\psi] \quad (5.36)$$

V can not be determined by method 2 but is more efficient to derive Q and U = most often used for stellar polarimetry

$$\left. \begin{aligned} I'(0^\circ) &= \frac{1}{2}(I + Q) & I'(45^\circ) &= \frac{1}{2}(I - Q) \\ I'(22.5^\circ) &= \frac{1}{2}(I + U) & I'(67.5^\circ) &= \frac{1}{2}(I - U) \end{aligned} \right\} \quad (5.37)$$



case of linear polarization (Method 2)

and solving for I, Q, and U gives

$$\left. \begin{aligned} Q &= I'(0^\circ) - I'(45^\circ) & U &= I'(22.5^\circ) - I'(67.5^\circ) \\ I &= I'(0^\circ) + I'(45^\circ) & I &= I'(22.5^\circ) + I'(67.5^\circ) \end{aligned} \right\} \quad (5.38)$$

5.3.3 Mueller matrices

Handles all four Stokes parameters simultaneously.

Each optical element represented by a 4x4 matrix.

Light represented by a 1x4 Stokes vector.

$$S' = M_n M_{n-1} \cdots M_2 M_1 S$$

$$M' = R(-\psi) M R(\psi)$$

$$R(\psi) = \begin{pmatrix} 1 & 0 & 0 & 0 \\ 0 & \cos 2\psi & \sin 2\psi & 0 \\ 0 & -\sin 2\psi & \cos 2\psi & 0 \\ 0 & 0 & 0 & 1 \end{pmatrix}$$

5.4 INTERFEROMETERS

- Collection method

Combining light from widely separated telescopes to overcome the diffraction limit of an individual telescope.

Applied in radio, but recently made the advancement of high-resolution interferometers in Opt+IR

- Detection method

Single-aperture telescopes with interferometer equipment for specific purposes. Several types of detection interferometers have been used for spectroscopy, such as the Fourier Transform Spectrometer(FTS) and the Fabry-Perot interferometer which is an imaging spectrometer.

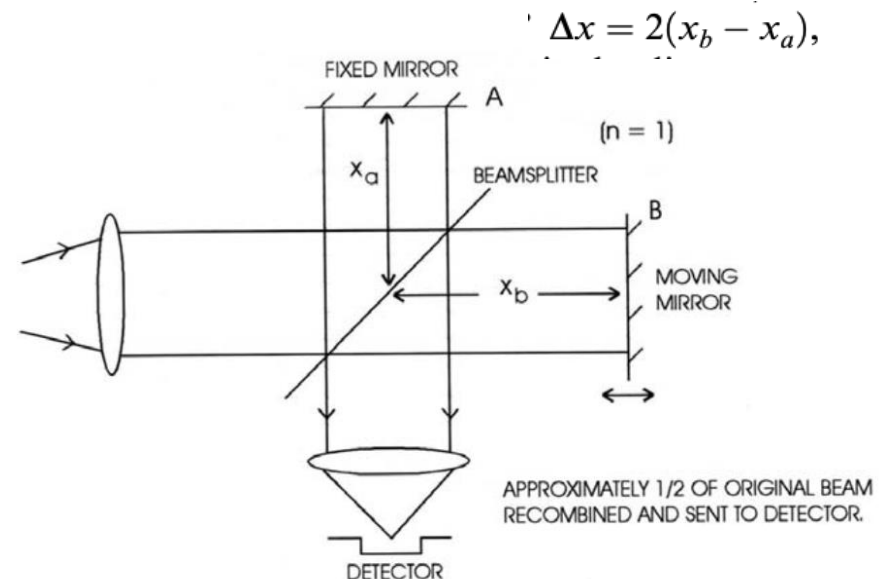
5.4.1 The Fourier Transform Spectrometer (FTS)

- Scanning Michelson interferometer with collimated light input.
- Measures intensity variations due to path difference between fixed and scanning mirrors.
- Fourier Transform: Converts interferogram to spectral information.
- Advantages
 - High Resolving Power: Example $4\Delta x_{\max}/\lambda$ and with $\Delta x_{\max} = 10$ cm
 - $R = 400,000$ at 1 μm wavelength with a 10 cm path difference.
 - High Signal-to-Noise Ratio: All light falls on the detector.
- Disadvantages

Time-Dependent Measurements: Atmospheric conditions may vary during measurement sequence.

$$T(k, \Delta x) = \frac{1}{2} [I + \cos(2k\Delta x)] \quad (5.41)$$

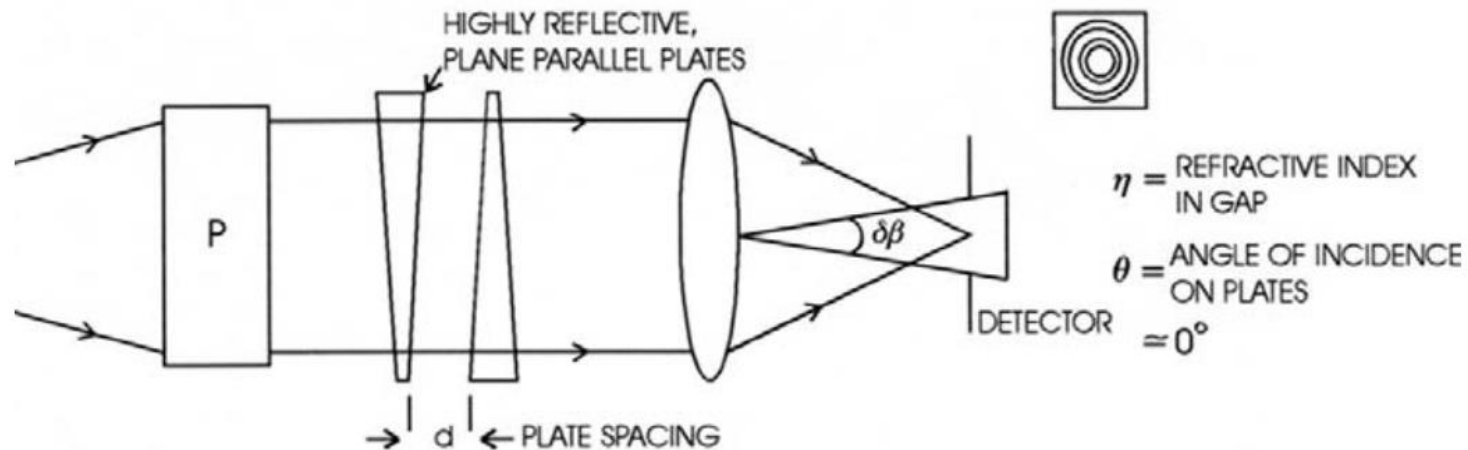
$$F(\Delta x) = c \int I(k) T(k, \Delta x) dk = \text{constant} + \frac{c}{2} \int I(k) \cos(2k\Delta x) dk \quad (5.42)$$



5.4.2 The Fabry-Perot etalon

The Fabry-Perot interferometer = imaging spectrometer which is formed by placing a device called an "etalon" (French *étalon*, meaning "measuring gauge" or "standard") in the collimated beam of a typical camera system.

- Two plane-parallel plates with highly reflective coatings.
- Wavelengths transmitted with maximum intensity follow the relation $m\lambda = 2nd\cos\theta$.
- Etalon in the collimated beam of a camera system.
- Set of concentric rings for monochromatic light.
- Pre-Filtering: Narrow band interference filter to ensure narrow light band.
- Resolving Power: Determined by the finesse of the etalon.



5.4.2 The Fabry-Perot etalon

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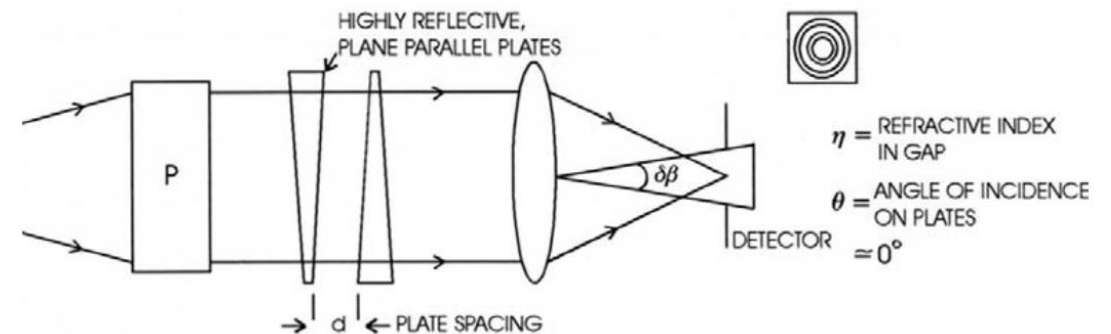
- Wavelengths transmitted with maximum intensity follow the relation $m\lambda = 2nd\cos\theta$.
- Etalon in the collimated beam of a camera system. Set of concentric rings for monochromatic light.
- Pre-Filtering: Narrow band interference filter to ensure narrow light band.
- Taurus Tunable Filter: Wide-field narrow-band imaging in the CCD range (370 nm to 1,000 nm).
- Finesse (F): Measure of plate quality and reflectance.
- Resolving Power (R): $R = \lambda/\Delta\lambda$, dependent on the finesse and spacing of the etalon plates.
- Charge shuffling synchronized to frequency (band) switching is used to suppresses systematic errors, enhancing image quality.

a circular aperture isolates the central order which has an angular diameter $\delta\beta = \sqrt{8/R}$ and the free spectral range is given by

$$\Delta\lambda_{FSP} = \frac{\lambda}{m} = \frac{\lambda^2}{2nd} \quad (5.44)$$

The resolving power ($R = \lambda/\delta\lambda$) is

$$R = \frac{2Fnd}{\lambda} \quad (5.45)$$



5.4.3 Interference filters

- Multi-layer thin-film devices operate with the same principle as Fabry-Perot etalon.
- Produces transmission maxima when wavefronts are in phase.
- Construction
 - Two quarter-wave stacks separated by a half-wave spacer.
 - Can have 3-4 layers for steeper band slopes and near square "tophat" profiles.
 - Continuous vacuum deposition run to create multiple layers.
 - Base width where transmission is 1% of peak is 1.9-2.2 times the FWHM.
 - Tilt-Scanning: Shifts center wavelength to the blue with increasing angle of incidence.

$$\lambda = \lambda_0 \sqrt{[1 - (n_o/n_e)^2 \sin^2 \varphi]} \quad (5.46)$$

