# AGNAGN Seminar Sec12-12.3

Misato Fujii

#### 12.1 Introduction

- · novae & supernovae: shells of gas are cast off from evolving star and returned to interstellar space
  - much more violent than planetary nebulae
    - velocity of expansion
      - nova shells:  $\sim 10^3$  km/s
      - supernova shells:  $\sim 10^{3-4}$  km/s
  - mass of nova shells are much smaller than planetary-nebula shells
    - planetary-nebula shells:  $\lesssim 10^{-4} M_{\odot}$
    - supernova shells:  $\geq 1M_{\odot}$
  - observed line spectra have general similarities
- heated and ionized gas tends to radiate more or less the same emission-line photons, regardless of the mechanisms
- →discuss each shells

#### 12.2 Nova Shells

- · novae: evolving star suddenly becomes much brighter, reaching  $L \gtrsim 10^4 L_{\odot}$ 
  - spectra $\rightarrow$ material leaves the star with  $v\sim10^3$  km/s
  - · emission-line profiles→indicate multiple "shells"
  - release  $E > 10^{45}$  erg/year, and gradually return to pre-outburst state ( $\approx 10^2 \, \mathrm{yr}$ )
- physical mechanism of novae
  - close binary stars (white dwarf & red dwarf or subgiant)
  - →overflow subgiant's Roche lobe and lose mass into white dwarf's lobe
  - →inflowing mass spirals into white dwarf as an accretion disk
  - →build-up of hydrogen-rich material on the surface of the white dwarf
  - →the temperature at its inner edge rises
  - →start explosive thermonuclear runaway
  - →energizes the nova outburst
  - ejecta: mixture of gas with a composition from the subgiant and material from the white dwarf
  - nova explosions repeat every  $\sim 10^4 10^5$  yr

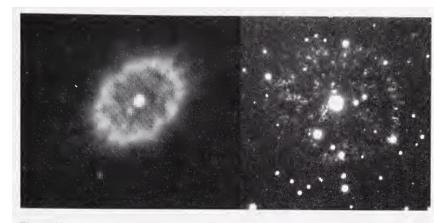


Figure 12.1

The shells around the novae DQ Her (left) and GK Per. DQ Her erupted in 1934 and its shell had a diameter of roughly 20" when this image was taken. GK Per erupted in 1901. (GK Per, WIYN Observatory)

# Spectrum

- initial stage
  - dense and optically thick from UV to infrared
  - form a photosphere similar to A-F supergiant with broad blue-shifted absorption lines
- · soon after the peak luminosity
  - · the density of the shell falls, the gas becomes optically thin
  - begins to show emission lines of HI and Hel
- · as the continuum weakens
  - the emission lines strengthen
  - typical nebular lines ([NII], [OIII], [NeIII]) start to appear and become stronger
    - initially, the density is above the critical of the nebular lines
      - they are faint
      - [NII] $\lambda$ 5755 and [OIII] $\lambda$ 4363 are relatively strong due to the high density
    - as the density falls below the critical densities
      - [NII] $\lambda$ 5755 and [OIII] $\lambda$ 4363 weaken with respect to [NII] $\lambda\lambda$ 6548, 6583 and [OIII] $\lambda\lambda$ 4959, 5007
- spectrum gradually changes to an almost pure nebular-type spectrum
  - · broad emission lines resulting from the high expansion velocity

## Growth of nova shell

- a few years after the outburst
  - · the ejecta becomes a small, faint nebulous shell surrounding the post-nova star
    - →expand at constant-velocity
    - →gradually become fainter
    - →eventually merge into and become part of ISM
- distance of a nova shell
- · determined by comparing the measured radial velocity of expansion with the proper motion (or angular velocity) of expansion
  - DQ Her (Nova Herculis 1934)
    - measured radial expansion velocity:  $320 \pm 20 \text{ km/s}$
    - the dimensions of the shell: 11" x 17"
    - $\rightarrow$  distance: 420 ± 100 pc
    - · uncertainty: the major or minor axis of the elliptical shell

# Spectra of shell

- · nova shells have quite unusual nebular spectra
  - many permitted lines of various stages of ionization of N and C
  - normal HI, Hel and Hell lines
  - a very few forbidden lines ([NII] $\lambda\lambda$ 6548, 6583 and [OII] $\lambda$ 3121)
- · reason of the nearly complete absence of forbidden lines from these optical spectra
  - the temperature is so low that they all are suppressed
    - thresholds for excitation  $\chi$  is much larger than thermal energies  $(\exp(-\chi/kT)\ll 1)$
    - · at low temperatures, only recombination lines are expected to be visible in the spectra of nebula
      - for a collisionally excited line near 6000Å, T < 3,500 K is required for  $\exp(-\chi/kT) < 10^{-3}$
- the strongest expected lines in the optical spectrum are H  $\alpha$  H  $\beta$  , Hel $\lambda$ 5876, and Hell $\lambda$ 4686 for HI and He
  - · observed in nova shells

### abundance of the elements in nova shells

- relative abundances of the elements in nova shells can be determined from the relative strengths of their emission lines
  - CII  $(2s^22p^2p^o 2s2p^2D \lambda 1335)$ : strong because of dielectronic recombination of  $C^{++}$  through  $C^{+}$
- observational data→nova shells are somewhat enriched in He, and greatly enriched in C, N, & O, compared to unevolved stars
  - ⇒significantly contribute to the present abundances of heavy elements in ISM
- · large abundances of C, N, and O
  - →explain the low temperature and the near absence of forbidden lines in the observed spectrum
    - even at low temperatures, the rate of collisionally excited far-infrared line radiation is very large pequilibrium temperature is quite small by higher radiative cooling rates

## Model of a nova shell

- photoionization model of a nova shell can be calculated by similar ways of HII regions and planetary nebulae
  - In a nova, the source of the ionizing radiation is the accretion disk around the white dwarf
  - the shell is not a simple sphere and the spectra of the minor and major axes are different
    - 1. genuine composition or density differences
    - 2. the more distant portions of the shell receive a lower flux
    - 3. the shape of the ionizing continuum may be different
    - best models (Table 12.1) make distinctions between the major and minor axes
      - →reproduces qualitatively all its main features
        - · quite different from the ordinary HII regions and planetary nebulae
- · the high-energy photon processes is taken into account to calculation
  - main source of excitation of [OII]λ3727
    - photoionization of  $\mathcal{O}^0$  by photons
    - $\leftrightarrow$  collisional excitation of  $O^+$  in typical HII regions and planetary nebulae
- main cooling source in nova shells: collisional excitation and line radiation from  $N^+$ ,  $N^{++}$ , and  $O^{++}$ 
  - · observation of these emission lines provide a direct measure of the major coolants

#### 12.3 The Crab Nebula

- supernovae: much higher luminosities than novae  $(M_{bol} \approx \sim -18 -20, L \approx 10^9 10^{10} L_{\odot})$ 
  - Type II: show hydrogen emission lines
    - · imply the presence of hydrogen in the outer envelope of the evolved star
    - end stages in the evolution of massive stars  $(M \gtrsim 8M_{\odot})$ 
      - · H burning, He burning, and further thermonuclear burning stages lead to heavy nuclei
      - →central core collapses to a neutron star or black hole, and a shell is expelled with high velocity
    - eject greater masses ( $\sim 10 M_{\odot}$ ) and lower luminosities ( $M_{bol} \approx -18$ )

# Classification of supernovae

- Type I: don't show hydrogen emission lines
  - →suggest that they originate from a source with little hydrogen
  - Type Ia: show Si lines
    - resulting from the thermonuclear destruction of white dwarf
      - accrete white dwarfs in binary systems
        - →grow to exceed the Chandrasekhar limit
        - →ignite a C (or He) detonation
    - progenitors or outputs are a homogeneous group
      - →used as standard candles to large distances
    - Type la eject lower masses ( $\sim 1 M_{\odot}$ ) and higher luminosities ( $M_{bol} \approx -19.6$ )
  - Type Ib: don't show Si lines, but have He lines
  - Type Ic: don't show Si nor He lines
- Type Ib and Ic: massive stars which have lost their hydrogen-rich outer envelop by stellar winds or the explosion that is similar to Type II

#### Crab Nebula

- · Crab Nebula (NGC 1952): the remnant of a supernova
  - show hydrogen emission lines with a relatively strong
    - →Type II supernova
    - the emission lines are concentrated in the filaments (Figure 12.2a)
    - the continuum arises in the amorphous gas (Figure 12.2b)
  - the radial velocity splitting at the center: ~2900 km/s
  - →expansion velocity: 1450 km/s
    - · much larger than in planetaries but considerably less than typical supernova
  - a pulsar very near the center of the Crab Nebula
    - →the neutron star remnant of the original pre-supernova star
  - an extremely strong radio source
    - · emission mechanism that produces both the radio-frequency and amorphous optical continuum
      - synchrotron emission

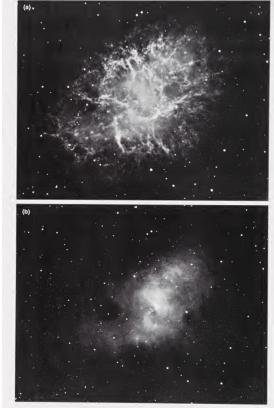


Figure 12.2

NGC 1952 (Crab Nebula). Both photos were taken with the Shane 3-m reflector. The upper photo (a) was taken with a red filter-plate combination transmitting chiefly Hα, [N II]. The lower photo (b) was taken with a yellow filter-plate combination transmitting chiefly continuum radiation. (ΘUC Regents/Lick Observatory)

# emission-line spectrum of the Crab Nebula

- · typical nebular lines (HI, HeI, HeII, [OII], [OIII], [NII], etc.)
- a wide range of ionization, including relatively strong [OI], [SII], and [NeV]
- [OII] $\lambda\lambda$ 3726, 3729 and [SII] $\lambda\lambda$ 6716, 6731 line ratios
  - $\rightarrow$ typical electron densities  $n_e \approx 10^3 \, \mathrm{cm}^{-3}$
- [OIII]  $(\lambda 4959 + \lambda 5007)/\lambda 4363 \rightarrow$  mean temperature T = 15,000 K
  - the corresponding ratio for [NII]  $\rightarrow$  T = 7,400 K
- · derived average abundances over the entire nebula
  - $\cdot [n(He^+) + n(He^{++})]/n(H^+) = 0.47$
  - $[n(0^+) + n(0^{++})]/n(H^+) = 10^{-3.5}$
  - $\cdot n(N^+)/n(H^+) = 10^{-4.0}$
  - ⇒ the material in the filaments is He-rich (result of nuclear processing)
    - abundances of N and O compared to H are approximately normal
      - $\cdot$  low compared to H + He

#### Structure of Crab Nebula filaments

- derived temperatures and abundances
  - ⇒photoionization is the main mechanism to the ionized gas in the filaments
  - the continuum of the amorphous region of NGC 1952 through the optical into the near UV
    - →can be fit with the X-ray flux from the Crab Nebula
  - optical continuum is strongly polarized
    - luminosity and polarization fit smoothly to the radio-frequency region
  - · optical continuum of the amorphous region of NGC 1952 is due to synchrotron emission
    - the radio continuum, and the UV and X-ray continua, too
- · magnetic field of the rotating neutron star gives energy to the gas near it
  - part of energy goes into expanding the nebula
  - · the rest into accelerating electrons that produce the synchrotron radiation
  - ⇒filaments: regarded as high-density regions photoionized by source of continuum radiation
    - direct photographs→synchrotron source is extended
    - · kinematic studies → filaments are in a thick shell surrounding the extended source

#### Model of Crab Nebula filaments

- the observed and calculated line spectra agree qualitatively (Table 12.4)
  - models reproduce the observed great strengths of  $[OI]\lambda 6300$ , 6364 and  $[SII]\lambda\lambda 6716$ , 6731
    - result of the relatively large fraction of high-energy photons in the photoionizing spectrum
      - high-energy photons have a small absorption cross section
        - →a significantly larger "transition legion" than HII regions
        - $\rightarrow$ in this transition region,  $H^0$ ,  $O^0$ ,  $H^+$ ,  $S^+$ , and electrons can all coexist
        - →collisionally excited [OI] and [SII] lines are emitted
- the calculated H  $\alpha$  /H  $\beta$  is significantly larger than the recombination value (2.85)
  - $\cdot$  collisional excitation of H  $\alpha$  from its ground level in this same transition zone
- · estimate total mass of the gas in the filament from luminosity and mean density
  - $\rightarrow \sim 1.5 M_{\odot}$ 
    - · may be larger if there are significant amounts of nearly neutral gas protected from ionizing radiation