

AGNAGN seminar

Sec.11.4-11.5

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11.4 Physical Conditions from X-ray Spectroscopy

X-ray spectrum

Classification of X-ray(based on detector technology):

- **Soft X-ray(<2keV)**
Emission lines from various element
- Hard X-ray(>2keV)
No emission line from abundant element, except Fe

Soft X-ray spectrum of NGC1068

(Emission lines and recombination continua are formed by the same physical process with IR and optical)

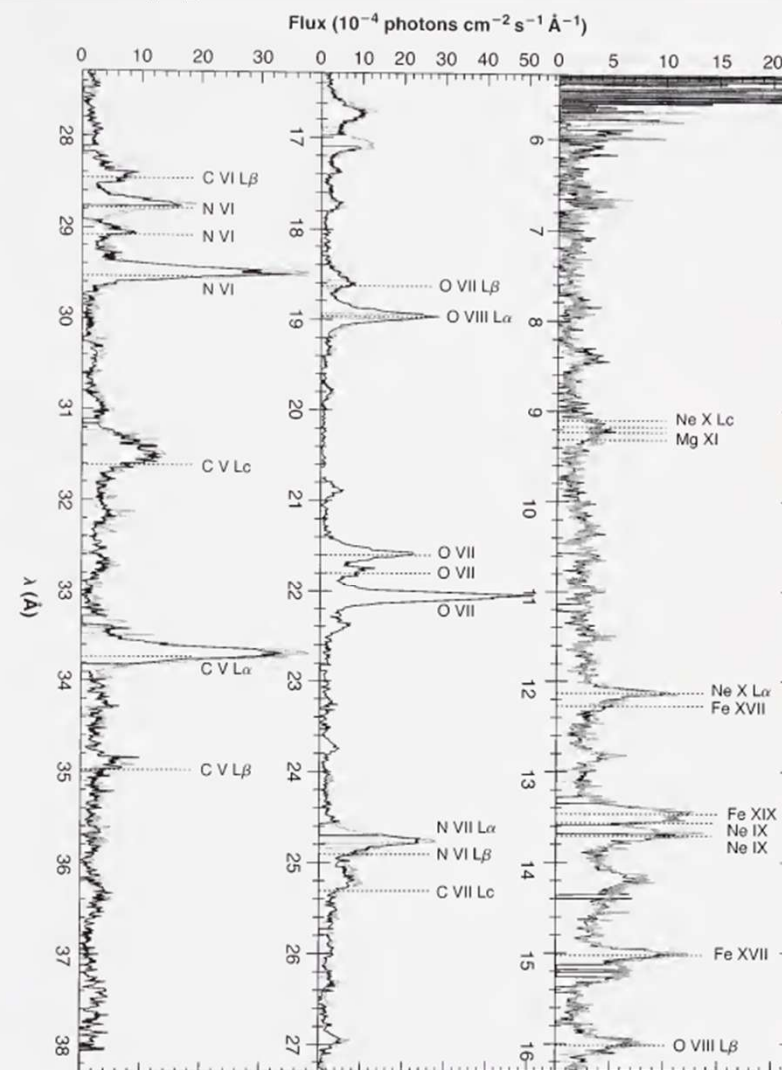


Figure 11.6

The soft X-ray spectrum of the Seyfert galaxy NGC 1068.

Representative emission lines: H-like ions

(Discuss in terms of **isoelectronic Sequences**)

- Ions with **1** electron: discuss the same as **H**
- Ions with **2** electrons: discuss the same as **He**
- **Transition probability: Z^4 Energy: Z^2**
(Z : atomic number)

Table 11.2

H-like $1s-2p$ wavelength and two-photon critical density

Ion	λ (Å)	$h\nu$ (keV)	Transition probability	Critical density
			$A(2\nu)$ (s^{-1})	$n_{crit}(2s)$ (cm^{-3})
C VI	33.74	0.366	3.8×10^5	5.1×10^{14}
N VII	24.78	0.499	9.7×10^5	1.7×10^{15}
O VIII	18.97	0.652	2.2×10^6	5.0×10^{15}
Si XIV	6.19	2.00	6.2×10^7	4.2×10^{17}
Fe XXVI	1.78	6.94	2.5×10^9	5.8×10^{19}

Gas temperature $<$ excitation energies

$\Rightarrow$ **Emission lines are produced by recombination**

$\Rightarrow$ ①(2/3) Populate 2p, and produce $L\gamma$ transition

②(1/3) **Populate 2s, and produce two-photon transition or $L\gamma$**

$\uparrow$ Gas density $\ll$ critical electron density n_{crit}

$\Rightarrow$ low density limit. two-photon transition is dominant

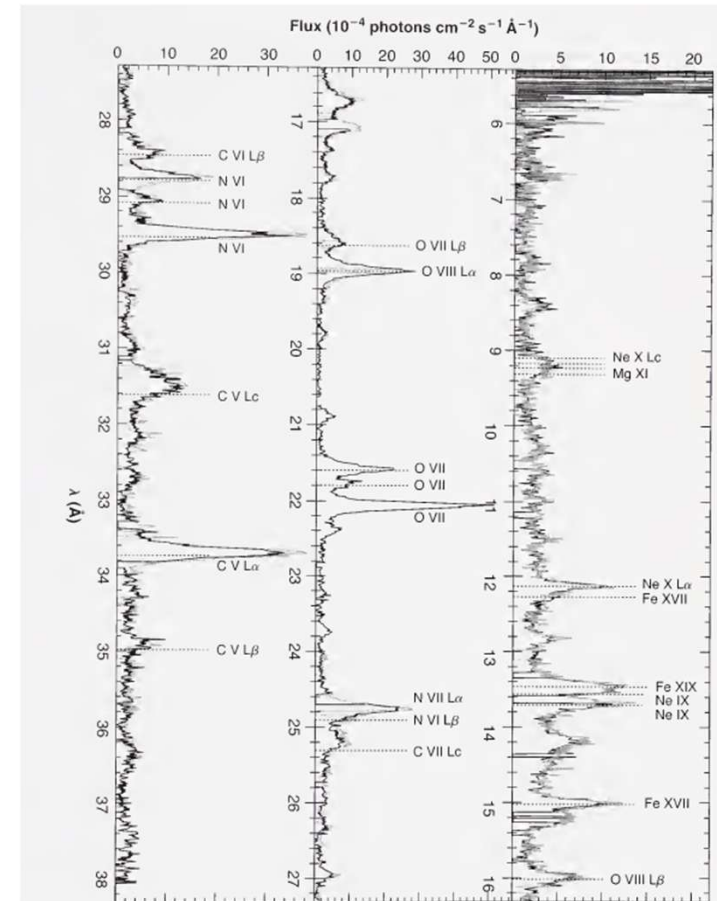


Figure 11.6

The soft X-ray spectrum of the Seyfert galaxy NGC 1068.

Higher Lyman recombination lines & Environment

① Case A

- **Optically thin**
- **continuum striking the clouds contains no radiation within the higher-n Lyman lines**
(appropriate for a cloud ionized by a hot star with strong Lyman absorption features)

② Case B

- **Optically thick**
- **Higher Ly photon $\Rightarrow$ Br photon + (Ly α or two-photon cont.)**

③ Case C

- **appropriate when the continuum is bright in the Lyman lines** (Similar environment with Case A, but contribution of continuum is strong)

Representative emission lines: He-like ions

Table 11.3

He-like $n = 2$ lines and transition probabilities

Ion	$1s^2\ ^1S_0 - 1s\ 2s\ ^3S_1$		$1s^2\ ^1S_0 - 1s\ 2p\ ^3P_1^o$		$1s^2\ ^1S_0 - 1s\ 2p\ ^1P_1^o$	
	$\lambda\ (\text{\AA})$	$A\ (\text{s}^{-1})$	$\lambda\ (\text{\AA})$	$A\ (\text{s}^{-1})$	$\lambda\ (\text{\AA})$	$A\ (\text{s}^{-1})$
C V	41.46	6.9×10^3	40.74	2.6×10^7	40.27	8.9×10^{11}
N VI	29.53	249	29.09	1.4×10^8	28.79	1.8×10^{12}
O VII	22.10	1000	21.81	5.5×10^8	21.60	3.3×10^{12}
Si XIII	6.743	3.4×10^5	6.687	1.5×10^{11}	6.645	3.8×10^{13}
Fe XXV	1.867	3.9×10^9	1.856	5.8×10^{13}	1.848	4.6×10^{14}

Two photon transition $\longleftrightarrow$ Transition from $n=2$

Transition probability is different
 $\Rightarrow$ relative intensities is sensitive to density

- Used for density diagnostics (range: $10^8\text{--}10^{12}\text{ cm}^{-3}$)
- Situation: Case B or C

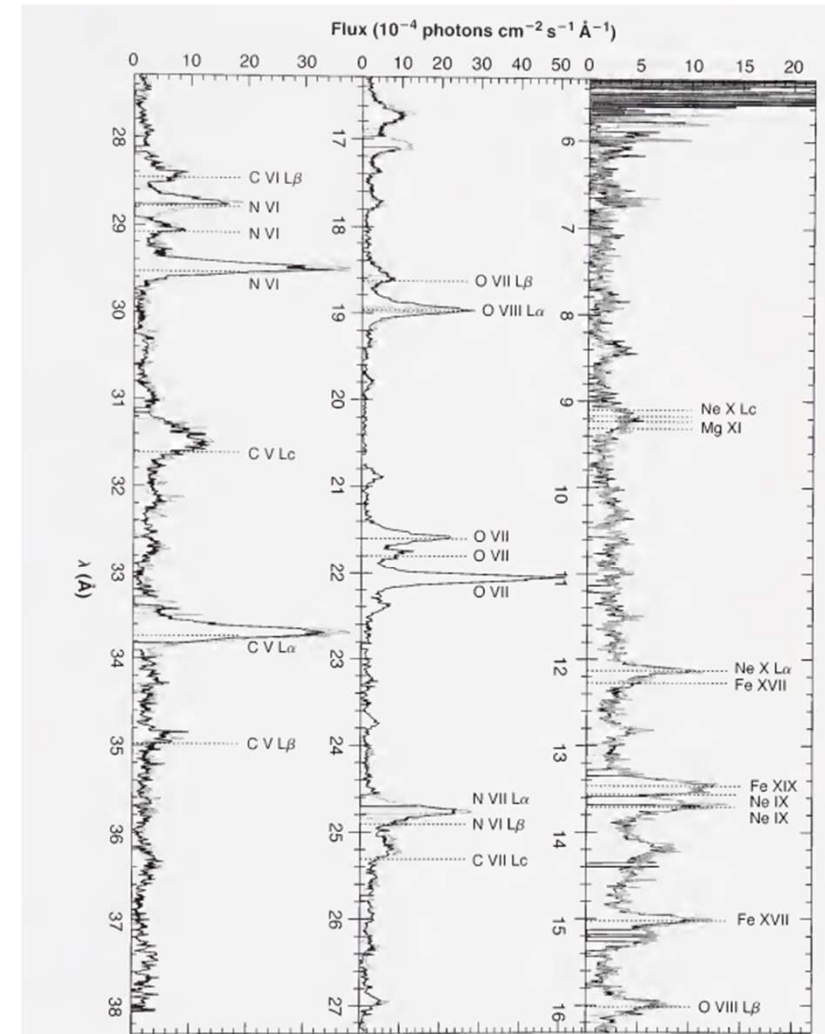


Figure 11.6

The soft X-ray spectrum of the Seyfert galaxy NGC 1068.

Characteristics of gas radiating X-ray

1. Highly ionized

- Many element have only 0~2 bound electrons
- Emit very little optical/IR, because of ①high energy transitions and ②low emissivity of optical/IR lines due to high temperature

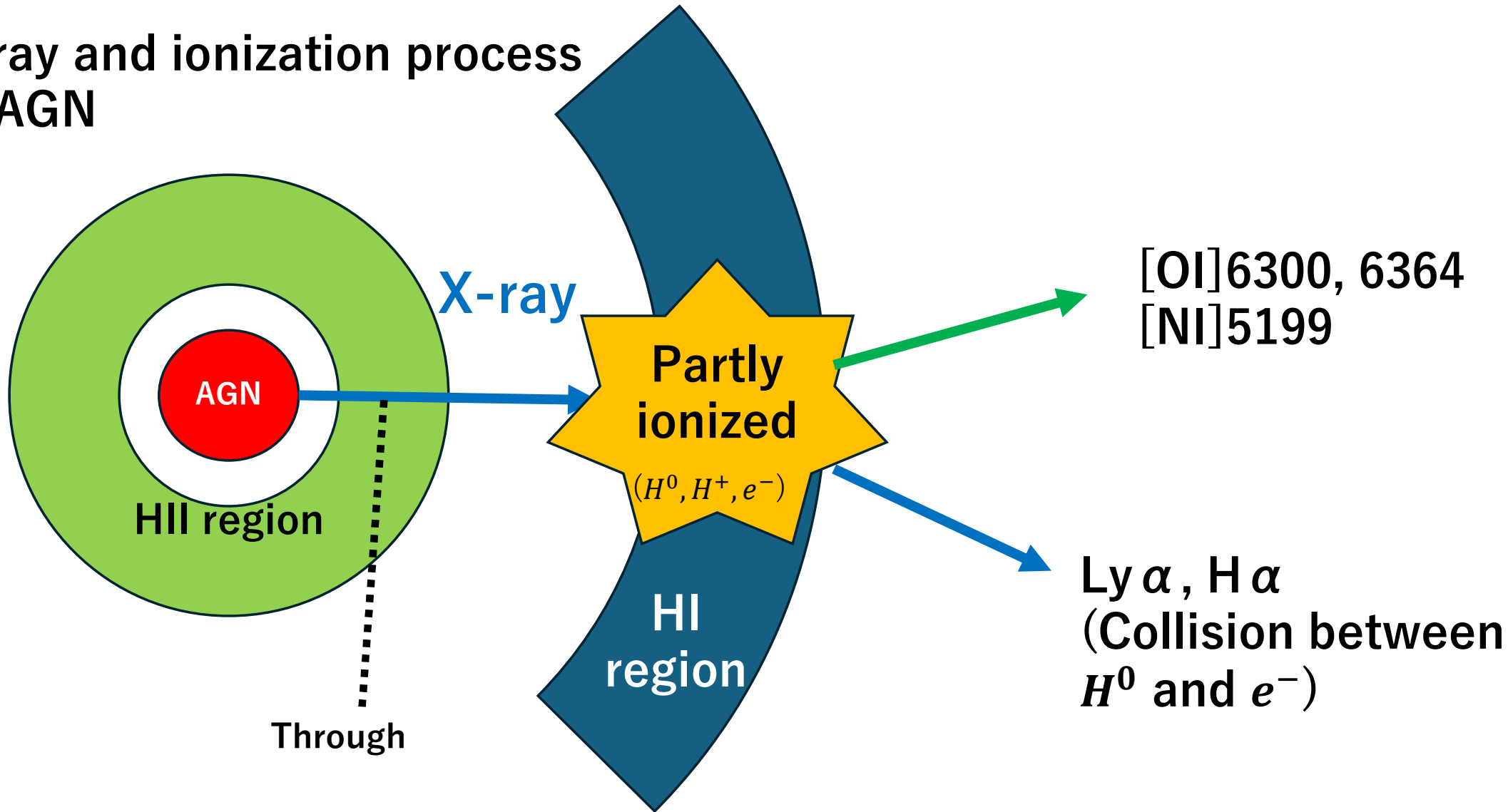
2. (Moderate ionized)

- Strong sources of optical/UV emission
- Produce fluorescent lines, but it needs large column densities(only see examples in AGN)(Sec.11.3)

Emission line in hard X-ray: Fe K α line (Sec. 11.2)

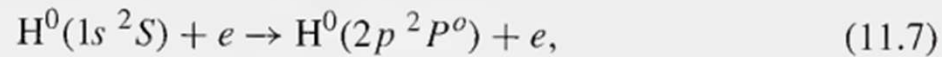
11.5 Collisional Excitation of H°

X-ray and ionization process
in AGN



Ly α collisional excitation (low density)

Reaction (direct excitation of the 2^2P^0):



Threshold: 10.2eV

Excitation cross section arises rapidly when (energy > threshold)

Emission coefficient (n=1 only, low density):

$$4\pi j_{L\alpha} = n_e n(\text{H}^0) q_{1^2S, 2^2P^0} h\nu_{L\alpha}, \quad (11.8)$$

Threshold and required collisional strength increase with n

⇒ **Contribution of collisional excitation to higher level will increase**
(about 3~10%)

Table 3.16

Effective collision strengths for H I

$T(\text{K})$	$1^2S, 2^2S$	$1^2S, 2^2P^0$	$1^2S, 3^2S$	$1^2S, 3^2P^0$	$1^2S, 3^2D$
10,000	0.29	0.51	0.066	0.12	0.063
15,000	0.32	0.60	0.071	0.13	0.068
20,000	0.35	0.69	0.077	0.14	0.073

⇒ consider higher n to derive (11.9)

Emission coefficient (n=1~3): **Contribution of higher n**

$$4\pi j_{L\alpha} = n_e n(\text{H}^0) (q_{1^2S, 2^2P} + q_{1^2S, 3^3S} + q_{1^2S, 3^2S} + q_{1^2S, 3^2D} + \cdots) h\nu_{L\alpha}. \quad (11.9)$$

Ly α emission coefficient(low density)

Equation for Recombination:

$$4\pi j_{L\alpha} = n_e n_p \alpha_{2\ 2p\ o}^{eff} h\nu_{L\alpha}$$

$$= n_e n_p (\alpha_B - \alpha_{2\ 2s}^{eff}) h\nu_{L\alpha} \quad (11.10)$$

Contribution of collisional excitation vs recombination:

Table 11.4

Ly α emission coefficients in partially ionized regions (all in $\text{erg cm}^3 \text{s}^{-1}$)

T (K)	Low density $n_e \ll 1.5 \times 10^4 \text{ cm}^{-3}$		High density $n_e \gg 1.5 \times 10^4 \text{ cm}^{-3}$	
	Collisional	Recombination	Collisional	Recombination
	$\frac{4\pi j_{L\alpha}}{n_e n(\text{H}^0)}$	$\frac{4\pi j_{L\alpha}}{n_e n_p}$	$\frac{4\pi j_{L\alpha}}{n_e n(\text{H}^0)}$	$\frac{4\pi j_{L\alpha}}{n_e n_p}$
10,000	2.56×10^{-24}	2.87×10^{-24}	4.46×10^{-24}	4.25×10^{-24}
12,500	2.62×10^{-23}	2.39×10^{-24}	4.51×10^{-23}	3.53×10^{-24}
15,000	2.23×10^{-22}	2.01×10^{-24}	2.10×10^{-22}	3.06×10^{-24}
20,000	8.89×10^{-22}	1.51×10^{-24}	1.46×10^{-21}	2.34×10^{-24}

High temperature/(High density)

⇒ **Collisional excitation is dominant**

Excitation of neutral H and contribution to Ly α (high density)

Critical density:

$$n_c = \frac{A_{2\ ^2S,1\ ^2S}}{q_{2\ ^2S,2\ ^2P}^p + q_{2\ ^2S,2\ ^2P}^e} \approx 1.5 \times 10^4 \text{ [cm}^{-3}\text{]} \quad (11.11)$$

Collisionally Emission coefficient(high density, n=1)

$$4\pi j_{L\alpha} = n_e n(H^0) (q_{1\ ^2S,2\ ^2S} + q_{1\ ^2S,2\ ^2P}) h\nu_{L\alpha} \quad (11.12) \quad \leftarrow (11.8)$$

Collisionally Emission coefficient(high density, n=1~3)

$$4\pi j_{L\alpha} = n_e n(H^0) \sum_{n=2}^3 \sum_{L=0}^{n-1} q_{1\ ^2S,n\ ^2S,n\ ^2L} h\nu_{L\alpha}. \quad (11.13) \quad \leftarrow (11.9)$$

Recombination coefficient(high density)

$$4\pi j_{L\alpha} = n_e n_p \alpha_B h\nu_{L\alpha} \quad (11.14) \quad \leftarrow (11.10)$$

Behavior are similar to “low density” case, but emission coefficient are larger by factors of approximately 1.5

H α & collisional/recombination

Under standard Case B condition, Collisional excitation to any of the levels 3^2S , 3^2P , or 3^2D leads to H α emission

Emission coefficient:

$$4\pi j_{H\alpha} = n_e n_p \sum_{L=0}^2 q_{1^2S,3^2L} h\nu_{H\alpha}. \quad (11.15)$$

- Strongly depends on temperature
- Not so important, comparison to Ly α (small contribution)

Table 11.5

H α emission coefficients in partly ionized regions (all in erg cm³ s⁻¹)

T (K)	Collisional $\frac{4\pi j_{H\alpha}}{n_e n(H^0)}$	Recombination $\frac{4\pi j_{H\alpha}}{n_e n_p}$
10,000	2.12×10^{-26}	3.54×10^{-25}
12,500	3.47×10^{-25}	2.89×10^{-25}
15,000	2.28×10^{-24}	2.46×10^{-25}
20,000	2.25×10^{-23}	1.81×10^{-25}