Mclean seminar sec.9.3-9.6

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- Pixel-to-pixel variations in sensitivity (QE) arise due to
 - Physical differences between pixels
 - optical attenuation effects
 - (e.g. microscopic dust particles on the surface of CCD)
- It is important to reduce pixel-to-pixel differences much further for astronomical observations.
 - Because such variations result in a "noisy" image
 - At a level corresponding to a few % of the sky brightness

- a common practice to overcome pix-to-pix variations
 - To observe the inside of the telescope dome
 - To place a huge white card on the dome
- the screens are so close that the telescope is completely out of focus
 - The field is uniformly illuminated (i.e. **flat**)
- Dome illumination is usually done with tungsten lamps.
 - Tungsten lamps do not mimic the spectrum of the night sky.



A) an image before flat process
B) Flat pattern
C) An image after flat processing (a is divided by b)

https://astro-dic.jp/flat-field-correction/

- For faint objects
 - It is the light of the sky that dominates.
 - It is better to try to use the sky itself as a flat-field
- For brighter objects
 - It is their own intrinsic color which matters.
 - inappropriate for using dome or sky
 - It is desirable to establish a set of narrow passbands for imaging
 - So as to limit the effect of color-dependent non-uniformity

- A simple arithmetic division pixel by pixel is required to remove pixel-to-pixel variations in sensitivity.
- *I_{FF}*: the uniform illumination of **the flat-field source** on a pixel in row i, column j
- η_{ij} : the quantum efficiency
- g : the conversion gain [electrons / DN]

$$\left(X_{ij}\right)_{FF} = \frac{1}{g}\eta_{ij}I_{FF}$$

• $(X_{ij})_{FF}$: the observed signal from that pixel in **DN**

• The mean signal in the flat-field is obtained by averaging X_{ij} over all the rows and columns : S_{FF}

$$S_{FF} = \frac{1}{g} \eta_M I_{FF}$$

• η_M : mean QE

• For the true image of the sky

$$X_{ij} = \frac{1}{g} \eta_{ij} I_{ij}$$

• To eliminate the position-dependent QE response, we form the ratio of the image.

$$\frac{X_{ij}}{\left(X_{ij}\right)_{FF}} = \frac{I_{ij}}{I_{FF}}$$

$$\frac{X_{ij}}{\left(X_{ij}\right)_{FF}} = \frac{I_{ij}}{I_{FF}}$$

$$S_{FF} = \frac{1}{g} \eta_M I_{FF}$$

- Finally, we rescale this ratio to the mean of the flat-field to give $\frac{X_{ij}}{\left(X_{ij}\right)_{FF}}S_{FF} = \left(\frac{\eta_M}{g}\right)I_{ij}$
- The flat-fielded, re-scaled image (left-hand) differs from the true image scene I_{ij} .

- Many flat-field exposures are averaged
 - to increase the accuracy of the flat-field itself
 - to remove from the flat-field various artifacts such as cosmic-ray events.
- The color of the flat-field should be as good a match as possible to that of the image scene
 - Because QE is a function of wavelength
- To detect the very weakest signals, more complex correction mean is necessary.
 - Obtaining a series of flat-fields at different exposure levels and determining the response of each pixel individually by means of a polynomial-fitting routine.

- For very high accuracy work or for observations on <u>objects</u> <u>fainter than the night itself</u>, **various systematic errors tend to dominate** over the expected random errors form photon arrival statistics.
- The first advance in counteracting low-level systematic errors
 - Drift scanning
 - Time delay and integration (TDI) (<u>Teledyne</u> page)
 - The CCD charge pattern is transferred slowly along columns, while the image from the telescope is physically scanned along with it **in precise synchronization**.

- The way to reduce systematic errors to a level of 0.1% of the night-sky background
 - Tony Tyson demonstrated
 - Theoretical limiting magnitude of 27th in a 6-hour exposure for a 4m class telescope
- The dominant sources of error
 - The mismatch in color between calibration flat-field and actual night-sky background
 - Interference fringing due to unblocked night-sky emissions (for thinned CCD)
- Generation of a master flat-field and sky frame from the **object frames** themselves removes systematic effects to better than 0.03% of night sky. (天体画像そのものからフラット画像を生成)





https://astro-dic.jp/dithering/

 This powerful technique (dithering) involves numerous observations of a piece of relatively blank sky with the telescope pointing to a slightly different position on the sky (displaced by 5-10 arcsec) for each exposure.

- Positions can be chosen randomly or in a simple pattern,
- But it is best not to repeat the pattern exactly.



https://astro-dic.jp/dithering/

- 1. The sequence of dithered exposures is examined
- 2. The frequency histogram examined
- 3. One signal value (or a small range) will turn out to **the pure night sky background**

(矢印のpixelに天体が載っていた り、載っていなかったりする。 画像枚数だけ、pixelカウントを 調べるとカウントの中央値が backgroundに相当)



Figure 9.8. (a) A raw CCD image with many defective pixels; (b) same image flattened by using "median sky flats" by shifting the images in a "dither" pattern. Credit: Harold Ables, U.S. Naval Observatory.

- One disadvantage of dithering
 - It will not work on object frames which are too crowded
 - a large galaxy
 - a large nebula
 - A centrally condensed cluster of stars
 - Unless much larger moves are made.
 - →A combination of dome flats and sky flats is often used.

- It is essential to normalize the various flat-fields **before** applying a median-filtering algorithm.
 - because a drift will affect the calculation of the median value.
- It is very important that any additive effects which do not vary with sky brightness should be removed **before scaling**.
- Additive effects
 - Electronics pattern noise (bias effect)
 - Charge skimming and trapping
 - Interference fringes
 - LED activity

- Multiplicative effect
 - QE variations across the sky
 - Transmission of optics and coatings
 - Thickness variations of thinned arrays and CCDs
- An appropriate steps for reduction and calibration of raw images
 - Subtract bias and bias structure (9.2)
 Subtract dark (9.2)
 Divide by flat-field (9.3)
 Subtract fringe frame (sky subtraction) (9.4)
 - 5. Interpolate(内挿する) over bad pixels : a bad pixel map is needed.
 - 6. Remove cosmic-ray events : identify non-star-like point sources.
 - 7. Registration of frames and median filtering

9.4 FRINGES AND SKY EMISSION



Figure 9.9. (a) A severe fringe pattern due to night-sky emission lines on a deep 4 m telescope exposure with a thinned, back-illuminated CCD. (b) The same field after processing to remove the fringes.

- Near-IR arrays can show significant fringe patterns in narrow-band work in the case of backside-illuminated CCD.
 - caused by OH emission lines
- The left figure shows 500s CCD exposure in a far-red band on a 4m telescope before and after fringe removal.

9.4 FRINGES AND SKY EMISSION

- Fringe removal can be performed by "adaptive modal filtering"
 - 1. Computes the **absolute difference between the mean and the median** of values associated with a pixel over all images in a set
 - 2. Rejects deviant values
 - until this difference falls below a certain value
 - or
 - until a maximum number of values have been rejected
- This technique fails if large or extended objects are seen.
- The introduction of deep-depletion CCDs and improved anti-reflection coatings has drastically reduced the problems.

9.5 LINEARITY

- It is usual that the **the output voltage** from a CCD **is proportional to the amount of light** falling on the CCD to very high accuracy.
- Linearity curves are usually derived by observing a constant source with various exposure times.
- Non-linear behavior from CCDs can occur if incorrect voltages are applied.
 - The output transistor is operating in its normal linear regime
 - Use the correct clock voltages to ensure the CCD pixel is fully inverted.
 - Use a CCD with MPP build in.

- Relative brightness ⇔ absolute amount of radiant energy reaching the Earth
 - A later method is needed if we are to understand the distribution of mass or energy in the Universe.
- Monochromatic flux (W m⁻² Hz⁻¹ / W m⁻² μm^{-1})
 - Integrating the specific intensity over the angular size of the source
- Magnitude
 - Relative monochromatic flux of a source
 - $m = m_0 2.5 \log F + 2.5 \log F_0$
 - m_0, F_0 : reference
 - There are different magnitude systems for different sets of <u>spectral bands</u>
 - Vega system / AB system

$$m = m_0 - 2.5 \log F + 2.5 \log F_0$$

- Vega system
 - Vega is assigned 0 magnitudes in every passband.
- AB system
 - F_0 is the same for all wavelengths and passbands.
 - $m_{\nu} = -2.5 log F_{\nu} 48.6$ (F is in frequency units)
 - $m_{\lambda} = -2.5 log F_{\lambda} 21.1$ (F is in wavelength units)
- Bolometric magnitudes
 - Gives a magnitude corresponding to the total flux integrated over all wavelengths
 - Zero point : $F_b = 2.52 \times 10^8 \text{ W m}^{-2}$
- Color indices
 - the difference between magnitudes at the two separate wavelengths (e.g. BV, UB)

- A UBV system is invented by Johnson and Morgan
- A modified UBVRI system



 Table 9.2. A summary of the major photometric systems.

Kron–Cous	ins System	Thuan–Gunn System		
Wavelength (Å)	Width (Å)	Wavelength (Å)	Width (Å)	
U 3,600	700	u 3,530	400	
B 4,400	1,000	v 3,980	400	
V 5,500	900	g 4,930	700	
R 6,500	1,000	r 6,550	900	
I 8,000	1,500	i 8,200	1,300	

Figure 9.10. Standard filter bandpasses used with CCDs: the Mould system.

- A red leak in the B filter
- The consequences of red leaks depend on the "color" of the illumination and the sensitivity of the detector at longer λ
- Calibration is needed.



Figure 9.11. The effect of (accidental) imperfect blocking is a "leak" of red photons to which the CCD is very sensitive. The consequence depends on the spectrum or color of the source; b and u are the balanced (filtered) and unbalanced artificial lamps illuminating the telescope dome, s is a typical solar spectrum, and t is the twilight sky which is quite blue.

- SDSS filters
 - wider bands to ensure high efficiency for faint-object detection
 - Covers the entire range where a CCD is sensitive (0.3-1.1 $\mu m)$



Figure 9.12. The Sloan Digital Sky Survey filter set.

Property	<i>u'</i>	g'	r'	i'	<i>z′</i>
$\lambda_{ m eff}$	355.1 nm	468.6 nm	616.5 nm	748.1 nm	893.1 nm
Width	56.0 nm	137.7 nm	137.1 nm	151.0 nm	94.0 nm
Limits	22.0	22.2	22.2	21.3	20.5

Table 9.3. Sloan Digital Sky Survey passbands and sensitivity limits.

- Two procedures for obtaining photometric information from CCD images
 - A) Aperture photometry
 - B) Profile fitting



an obtained image

- <u>Profile fitting</u> : Point Spread function (PSF) fitting
 - Modeling the image rather than summing over the image.
- Gaussian profile

$$I(r) = I_0 e^{-r^2/2\sigma^2}$$

- 1. Programs like DAOPHOT (Stetson, 1987) will identify the bright stars.
- 2. Deduce their Gaussian profiles and subtract those profiles away
- 3. Thereby revealing fainter stars.
- a convenient term

$$FWHM = 2.35\sigma$$

- Four important issues to be considered for fitting
 - 1. passband mismatch of the filters, including narrow-band filters
 - 2. Red/infrared "leaks" in the filter which complicate flat-fielding
 - 3. The finite opening and closing times of electromechanical shutters
 - 4. Changes in the atmospheric attenuation or "airmass" for long on-chip integrations
- 1 and 2 need to be eliminated by design
- shutter timing (3) errors are mainly relevant for large iris-type shutters that require a finite time (δt) to close,
 - The pixels at the center experience a longer exposure time.
 - For large arrays, the error could be ${\sim}0.1{\rm s}$
- Bonn shutter alleviates the problem.
 - <u>https://www.youtube.com/watch?v=d2dZQkqWIOQ</u>



(https://www.printables.com/model/566642-iris-shu

- Photometric values must be compared or calibrated against wellmeasured "standard" sources.
- Photometric standard stars must be observed over a wide range of airmass. $airmass(X) = \sec(90^\circ - \zeta) = \frac{1}{\cos(90^\circ - \zeta)}$



- the observed mag at $\zeta : m(\zeta)$
- The difference between the true mag and mag which would be observed at $\zeta:\alpha$

$$m = m(\zeta) - \alpha_{\lambda} \sec(\zeta)$$

• True mag and α_{λ} can be calculated by plotting the $m(\zeta)$ against airmass $(X = \sec(\zeta))$.



• The instrumental magnitudes

$$M = -2.5\log(counts/s)$$

- The parameters to be determined
 - the "zeropoint"
 - the "color equation" relating the CCD photometric system to the older photoelectric systems
 - the "extinction" factor or light-loss through the Earth's atmosphere per unit airmass $m = -2.5 \log(counts/s) \alpha * (airmass) + \beta * (color) + ZP$
- X (airmass) is changing for a long exposure.
 - DAOPHOT package

$$\bar{X} = \frac{X_0 + 4X_{1/2} + X_1}{6} + O(e)$$

- At the beginning X_0 , midpoint X_1/2 , the end X_1
- O(e) is a small error of about 1part in 10,000