
Dark Supernova Remnant

Yoshiaki Sofue¹

¹Institute of Astronomy, The University of Tokyo, 2-21-1 Mitaka, Tokyo 181-8588, Japan

*E-mail: sofue@ioa.s.u-tokyo.ac.jp

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Abstract

An almost perfect round hole of CO-line emission with a diameter of 3.7 pc was found in a molecular cloud (MC) centered on G35.75-0.25 ($l = 35^\circ.75, b = -0^\circ.25$) at radial velocity of 28 km s⁻¹. The hole is quiet in radio continuum emission, unlike the usual supernova remnants (SNR), and the molecular edge is only weakly visible in 8 and 24 μ m dust emissions. The hole may be either a fully evolved molecular bubble around a young stellar object (YSO), or a relic of a radio-quiet SNR that has already stopped expansion after rapid evolution in the dense MC as a buried SNR. Because G35.75 exhibits quite different properties from YSO-driven bubbles of the same size, we prefer the latter interpretation. Existence of such a "dark" SNR would affect the estimation of the supernova rate, and therefore the star formation history in the Galaxy.

Key words: galaxies: individual (Milky Way) — ISM: CO line — ISM: molecular clouds — ISM: supernova remnant

1 Introduction

Supernovae (SN) exploded in the Galactic disc with the mean gas density of $\sim 1 \text{ H cm}^{-3}$ expand to form supernova remnants (SNR), which emit thermal and non-thermal radiations due to the strong shock wave during their life time of $\sim 10^4$ y, and they finally merge into the interstellar medium (ISM) at radii exceeding a few tens of pc after $\sim 10^5$ y (Chevalier 1977; Raymond 1984; Weiler et al. 1988; Inoue et al. 2012). Such SNR are usually observed as extended objects, often having shell structures, by radio, X-ray, and optical measurements, and about 300 SNRs are currently catalogued (Green 2009).

Because of the short life time of high-mass stars responsible for SN explosions, SNRs are supposed to be located near to their birth places and adjacent MCs (Wheeler et al. 1980; Cox et al. 1999; Seta et al. 2004). SNe exploded in MCs evolve as buried SNRs quite differently from those exploded in the inter cloud space, diminishing in a short life time with a small radius (Shull 1980; Wheeler et al. 1980; Lucas et al. 2020). However, no observational evidence has been obtained yet of such 'buried' SNRs, mainly because of their compactness and short life time. This im-

plies that the current estimation of the SN rate from the catalogued SNRs in the Galaxy would have been significantly under-estimated.

In spite of the short life time of the buried SNR, the molecular cavity left after the rapid SNR evolution can survive for much longer time on the order of $t \sim r/\sigma_v \sim$ a few pc / a few km s⁻¹ $\sim 10^6$ y, which would be detectable as a hole in the MC by molecular line observations. Here, r is the radius of the shell, and σ_v is turbulent velocity in the cloud.

In this Letter, we report the discovery of an almost perfect round-shaped hole in a medium sized MC during the search for the sites of SNR-MC interaction (Sofue, et al. 2020), using the FUGIN (Four-receiver unbiased galactic imaging survey with the Nobeyama 45-m telescope) survey data in the ¹²CO, ¹³CO, and C¹⁸O ($J = 1 - 0$) line emissions (Umemoto et al. 2017). The full beam width at half maximum of the telescope was 15'', the velocity resolution was 1.2 km s⁻¹. The final 3D FITS cube data had an effective beam size of 20'', rms noise level of ~ 1 K, and a pixel size of $(\Delta l, \Delta b, \Delta v_{\text{lsr}}) = (8''.5, 8''.5, 0.65 \text{ km s}^{-1})$.

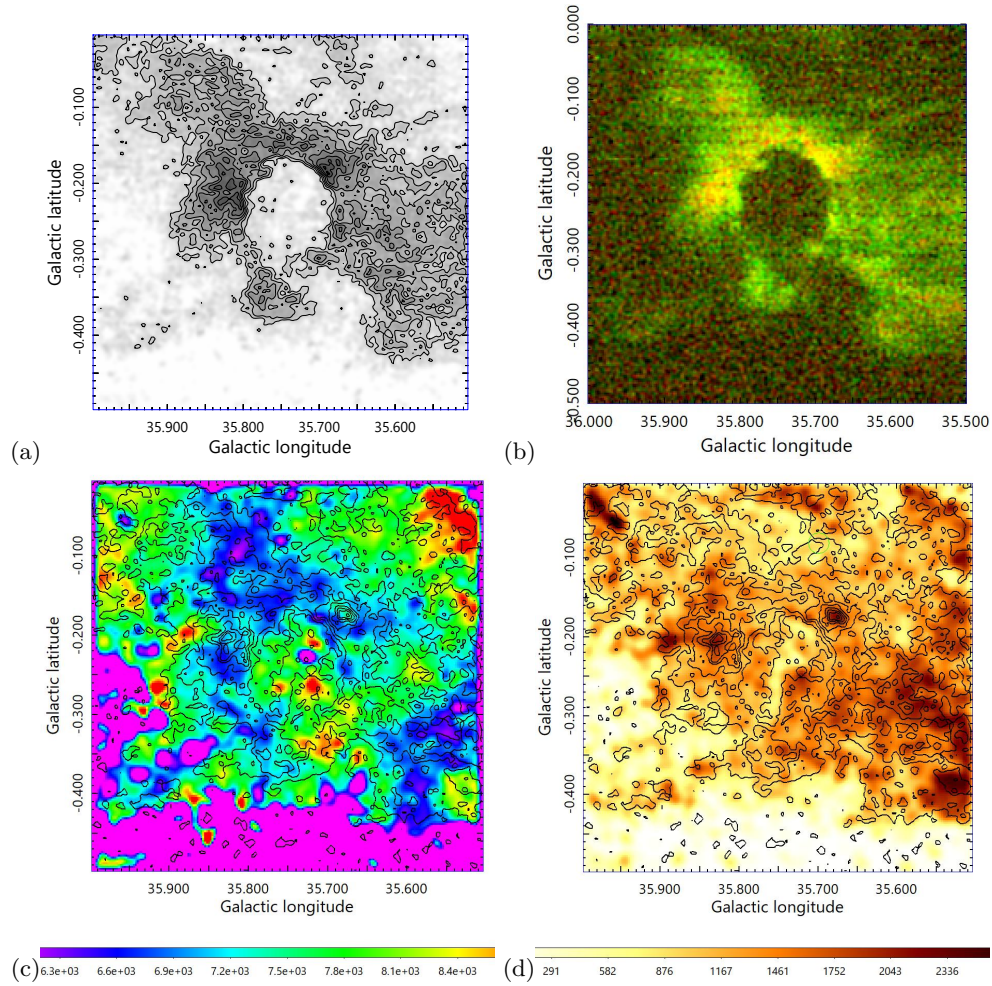


Fig. 1. (a) Integrated intensity of the ^{12}CO line from 24 to 31 km s^{-1} around the peculiar CO hole G35.75-0.25+28 km s^{-1} . Contours start from 10 with increment 5 K km s^{-1} . (b) Composite color map of the brightness temperature T_B of the ^{12}CO line at 27.775 km s^{-1} in green and ^{13}CO in red. (c) Moment 1 and (d) 2 maps showing velocity field (color code in m s^{-1}) and velocity dispersion superposed by intensity contours.

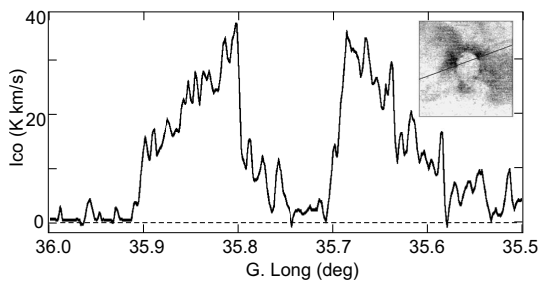


Fig. 2. Cross section of the ^{12}CO T_B across the line as inserted, showing that the hole's center is at the same level as the surrounding region. Such low brightness is explained only if the hole is open along the line of sight.

2 Peculiar CO hole G35.75-0.25

Figure 1 shows the distribution of integrated ^{12}CO intensity around the center velocity of 28 km s^{-1} centered on G35.75-0.25 ($l = 35^\circ.75, b = -0^\circ.25$), which shows a round-shaped hole of CO emission in a molecular cloud. Figure

2 shows the cross section of intensity across the hole. The figures show an almost perfect, round-shaped hole penetrating the molecular cloud with a steep edge revealed by the sharp intensity cut off. The surrounding molecular gas density is highest near the hole's edge, and decreases outward. Accordingly, the ^{13}CO emission in red has the maximum at the highest density regions with extended ^{12}CO envelope in green. The cross section indicates that the hole's edge is not compressed against the surrounding cloud, but is sharply cut from a supposed increasing density profile toward the center. This suggests that the hole was created in the cloud by dissociating the molecules, but not by sweeping the interior gas to the edge.

From the radial velocity at 28 km s^{-1} of the CO emission, the kinematic distance is estimated to be 1.8 ± 0.3 kpc for near-side solution, and 11.1 ± 0.3 for far solution using the newest rotation curve of the Milky Way (Sofue 2020). Because the molecular cloud

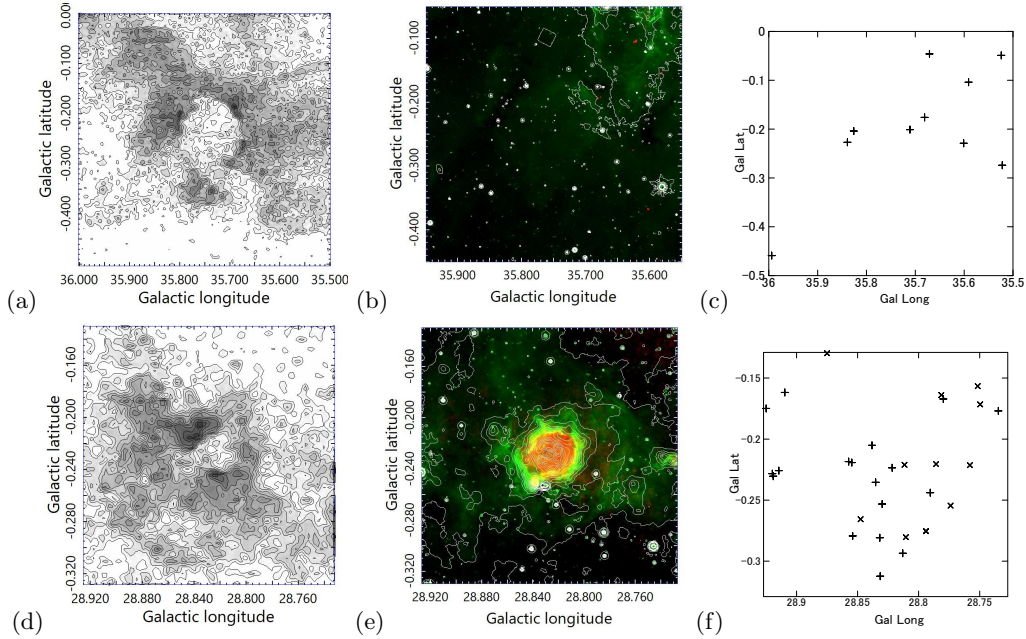


Fig. 4. Comparison with a Spitzer bubble N49 of about the same size. (a) ^{12}CO channel map of G35.75 at 28.8 km s^{-1} by contours (start 5 K, increment 2.5 K) (b) Color-coded map of 20-cm radio in red (1 to 3 mJy/beam), $8 \mu\text{m}$ in green (40 to 200 mJy/str), and $24 \mu\text{m}$ by white contours (1 to 1000 mJy/str by 0.1 dex log. increment). (c) Distribution of YSOs from SIMBAD marked by crosses (AGAL, 2MASS) (<http://simbad.u-strasbg.fr/simbad/>). (d),(e),(f) Same, but for a Spitzer bubble N49 at 85.2 km s^{-1} with the same intensity scales (Deharveng et al. 2010) with CO at 86.2 km s^{-1} , whose linear size is $\sim 4 \text{ pc}$, about same as G35.75. Note the quiet G35.75 in radio and infrared, and much higher YSO density in N49.

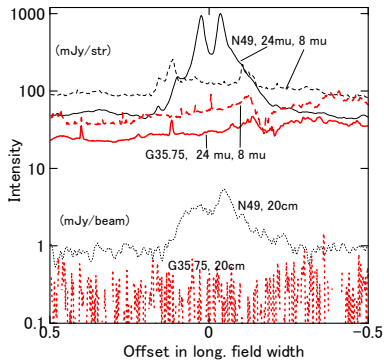


Fig. 5. Cross sections of the intensities across the centers of G35.75 and N49 from figure 4 relative to the field width in longitude. Note one to two orders of magnitude lower intensities of G35.75 in radio and infrared, in spite of comparable ^{12}CO brightness.

The evolution of a SNR in its evolved phase can be approximated by an adiabatic shock wave. Then, the expansion velocity v and radius r of the shock front is approximately related to the input energy E_0 and gas density ρ_0 as $E_0 \sim \frac{1}{2} \left(\frac{4\pi}{3} r^3 \rho_0 \right) v^2$. Remembering $v = dr/dt$, we have a differential equation for r as a function of time t , and the solution yields $v \sim 0.7(E_0/\rho_0)^{1/2} r^{-3/2}$, $v \sim 1.2(E_0/\rho_0)^{1/5} t^{-3/5}$, and $r \sim 0.5(E_0/\rho_0)^{1/5} t^{2/5}$. This leads to $v \sim 20 \text{ km s}^{-1}$ at $r \sim 2 \text{ pc}$ and $t \sim 5 \times 10^5 \text{ yr}$ for $n_{\text{H}_2} = 10^4 \text{ cm}^{-3}$ ($\rho_0 = \mu \times 10^4 \text{ H cm}^{-3}$, with $\mu = 2.8$).

If the cloud is large enough compared to the shell size, the SNR dies, being kept buried in the cloud, leaving a dark cavity, as illustrated in figure 6. However, if the cloud is flattened, the shock front is elongated in the polar direction, and finally open to the inter-cloud space. Accordingly, the internal pressure escapes to the polar direction, and the expansion in the equatorial plane will be decelerated more rapidly than in the spherical expansion.

Such a hole seen from the pole will appear as a round cavity having brightness as low as the surrounding background, as indeed observed (figure 2). Since the SNR's energy is almost exhausted at this stage, no particular kinematic disturbance to increase the velocity dispersion is observed in the cloud (figures 1(d), 3(c)). Even if the gas is expanding at the velocity as estimated above, since the motion is in the plane of the sky perpendicular to the line of sight, the velocity width at the edge will be not much increased compared to the ambient velocity dispersion, as indeed observed (figure 3(c)). This makes a contrast to the high velocity wing observed in a molecular cloud interacting with the bright SNR such as W44, which is expanding at $\sim 100 \text{ km s}^{-1}$ into the inter-cloud space with much lower density (Seta et al. 2004).

3D evolution of the shock front can be simulated by applying the radial adiabatic shock method (Sakashita 1971; Sofue 2019), as shown in figure 6, where a point

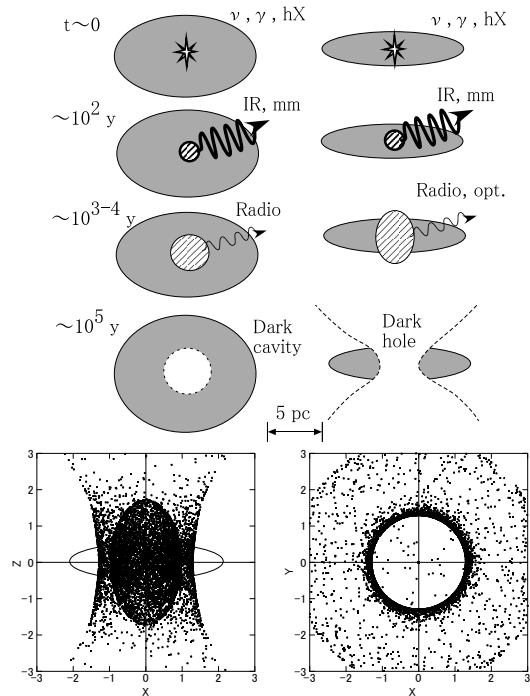


Fig. 6. (Top) Schematic evolution of a buried SNR in an extended molecular cloud (left panel) and in a flat cloud finally broken in the polar direction (right panel). (Bottom) 3D shock front in a flat spheroidal cloud by explosion at the center. Left and right panels show side and top views, respectively. Scale is arbitrary, and see the estimation in the text for the real scale and time.

explosion takes place at the center of a cloud having a Gaussian density profiles of radius 2 and thickness 0.5 in arbitrary units. As the time elapses, the shock front expands into the polar direction, and breaks the cloud to create a hole, mimicking the observed molecular hole having a clear-cut inner edge (figure 2).

4 Summary

A round-shaped hole of CO-line emission with a diameter of 3.7 pc was found in a molecular cloud centered on G35.75-0.25 at $v_{\text{lsr}} = 28 \text{ km s}^{-1}$, whose center is almost empty in radio and infrared emissions. The hole is interpreted as a relic of a fully evolved SNR that has already stopped expanding after its rapid evolution in the dense molecular gas as a buried SNR. Detection of such a "dark" SNR gives evidence for the existence of buried SNRs in the Galactic plane, probably every where. This would imply that the supernova rate in the Milky Way based on optical, radio and/or X-ray bright SNRs would have been significantly underestimated.

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